

AD-753 302

RECURRENCE FORMULA FOR THE VENEZIANO
MODEL N-POINT FUNCTIONS

Koichi Mano

Air Force Cambridge Research Laboratories
L. G. Hanscom Field, Massachusetts

1 May 1972

DISTRIBUTED BY:

NTIS

National Technical Information Service
U. S. DEPARTMENT OF COMMERCE
5285 Port Royal Road, Springfield Va, 22151

AD-753302

DOCUMENT CONTROL DATA - R&D		
(Security classification of title, body of abstract and indexing annotation must be entered when the overall report is classified)		
1. ORIGINATING ACTIVITY (Corporate author) Air Force Cambridge Research Laboratories (LZN) L.G. Hanscom Field Bedford, Massachusetts 01730		2a. REPORT SECURITY CLASSIFICATION Unclassified
		2b. GROUP
3. REPORT TITLE RECURRENCE FORMULA FOR THE VENEZIANO MODEL N-POINT FUNCTIONS		
4. DESCRIPTIVE NOTES (Type of report and inclusive dates) Scientific. Interim.		
5. AUTHOR(S) (First name, middle initial, last name) Koichi Mano		
6. REPORT DATE 13 December 1972	7a. TOTAL NO. OF PAGES 23	7b. NO. OF REFS 7
8a. CONTRACT OR GRANT NO.	9a. ORIGINATOR'S REPORT NUMBER(S) AFCRL-72-0723	
a. PROJECT, TASK, WORK UNIT NOS. 86820201		
c. DOD ELEMENT 62101F		
d. DOD SUBELEMENT 681300	9b. OTHER REPORT NO(S) (Any other numbers that may be assigned this report)	
10. DISTRIBUTION STATEMENT Approved for public release; distribution unlimited.		
11. SUPPLEMENTARY NOTES Reprinted from J. Math. Phys., Vol. 13, No. 10, October 1972.		12. SPONSORING MILITARY ACTIVITY Air Force Cambridge Research Laboratories (LZN) L.G. Hanscom Field Bedford, Massachusetts 01730
13. ABSTRACT A recurrence formula is derived for a function which reduces to the Veneziano model $(n + 3)$ -point function. It is shown that the formula is equivalent to, but is more self-contained than, the Hopkinson and Plahte formula in that it does not require the prescription for the parameters involved.		
Reproduced by NATIONAL TECHNICAL INFORMATION SERVICE U.S. Department of Commerce Springfield VA 22151		
KEYWORDS: Veneziano-type amplitudes		

Recurrence Formula for the Veneziano Model N -Point Functions

Koichi Mano

Air Force Cambridge Research Laboratories, Bedford, Massachusetts 01730
(Received 1 May 1972)

A recurrence formula is derived for a function which reduces to the Veneziano model $(n+3)$ -point function. It is shown that the formula is equivalent to, but is more self-contained than, the Hopkinson and Plafie formula in that it does not require the prescription for the parameters involved.

The extension to the n -point functions of the Veneziano's four-point function was accomplished either by generalizing the integral representation for the beta function which comprises the essential ingredient of the four-point function^{1,2} or by generating a recurrence formula for the n -point function.³ In the latter approach, the authors attempted to justify the formula for arbitrary value of n after showing, through introduction of the integral representation for the beta function, that the formula produces the already known integral expressions for the cases of $n = 5, 6$, and 7 . The recurrence formula which has apparently been discovered on a heuristic basis is not necessarily very transparent, as the authors themselves admit it, especially in connection with the definition of their variables x_{ij} .

Recently it was pointed out that the generalized Veneziano amplitudes may be regarded as the boundary values of a class of generalized hypergeometric functions that are Radon transforms of products of linear forms.⁴ In a work by the present author which shows that the amplitudes possess a structure similar to that of the Lauricella's hypergeometric functions,⁵ he has made an iterative use of a recurrence formula for the amplitude.⁶

The purpose of this note is to point out that the author's recurrence formula obtains itself in a quite natural manner such that for a special choice of the variables, it reproduces the formula proposed in Ref. 3 without requiring any prescription for the parameters involved therein.

To begin our discussion, let us consider a function V_n of variables w_{ij} defined as below:

$$\begin{aligned} V_n(\alpha_{01}, \alpha_{02}, \dots, \alpha_{0n}; \alpha_{11}, \alpha_{12}, \dots, \alpha_{1n}; \\ \alpha_{21}, \alpha_{22}, \dots, \alpha_{2,n-1}; \alpha_{31}, \dots, \alpha_{3,n-2}; \dots; \\ \alpha_{n-1,1}, \alpha_{n-1,2}; \alpha_{n1} | w_{02}, w_{03}, \dots, w_{0n}; \\ w_{13}, \dots, w_{1n}; \dots; w_{n-2,n}) \\ = \int_0^1 \dots \int_0^1 \prod_{i=1}^n \left\{ du_i u_i^{\alpha_{0i}-1} (1-u_i)^{\alpha_{1i}-1} \right. \\ \left. \times \prod_{k=2}^i \left[1 - u_i \left(w_{i-k,i} \prod_{j=0}^{k-2} u_{i-j-1} \right) \right]^{\alpha_{k,i-k+1}} \right\}. \end{aligned} \quad (1)$$

When the parameters α_{ij} are regarded as functions of the momenta p_i , $i = 1, 2, \dots, n+3$, of the external particles, it will be seen that V_n for $w_{ij} = 1$ can readily be related to the known integral representation for the $(n+3)$ -point function.⁷

In carrying out the multiple integrations in Eq. (1), use has been made in Ref. 6 of the following recurrence formula:

$$\begin{aligned} V_n(\alpha_{01}, \dots, \alpha_{0n}; \alpha_{11}, \dots, \alpha_{1n}; \dots; \alpha_{n1} \\ | w_{02}, \dots, w_{0n}; w_{13}, \dots, w_{1n}; \dots; w_{n-2,n}) \\ = B(\alpha_{0n}, \alpha_{1n}) \sum \frac{(\alpha_{0n}, \beta_{n,n-1})}{(\alpha_{0n} + \alpha_{1n}, \beta_{n,n-1})} \end{aligned}$$

$$\begin{aligned} \times \frac{(-\alpha_{2,n-1}, r_{2,n-1}) \dots (-\alpha_{n1}, r_{n1})}{(1, r_{2,n-1}) \dots (1, r_{n1})} \\ \times (w_{n-2,n})^{r_{2,n-1}} (w_{n-3,n})^{r_{3,n-2}} \dots (w_{0n})^{r_{n1}} \\ \times V_{n-1}(\alpha_{01} + \beta_{n1}, \dots, \alpha_{0,n-1} + \beta_{n,n-1}; \alpha_{11}, \dots, \\ \alpha_{1,n-1}; \dots; \alpha_{n-1,1} | w_{02}, \dots, w_{0,n-1}; w_{13}, \dots, w_{1,n-1}; \\ \dots; w_{n-3,n-1}). \end{aligned} \quad (2)$$

In Eq. (2) the summation is over the integers between 0 and ∞ of the indexes r_{ij} , B stands for the beta function, (α, r) under the summation symbol is written for $\Gamma(\alpha + r)/\Gamma(\alpha)$, and $\beta_{p,q}$ are given by

$$\beta_{p,q} = \sum_{k=p-q+1}^p r_{k,p-k+1} \quad \text{for } \begin{cases} p=2,3,\dots,n \\ q=1,2,\dots,p-1 \end{cases}. \quad (3)$$

Our task in what follows is to show that when $w_{ij} = 1$, Eq. (2) reduces to the recurrence formula for the N -point function $B_N(x)$ in Ref. 3 as given below:

$$\begin{aligned} B_N(x) = \sum_{k_{i,N-1}=0}^{\infty} \left[\prod_{i=2}^{N-3} (-1)^{k_{i,N-1}} \binom{z_{i,N-1}}{k_{i,N-1}} \right] \\ \times B_4(x_{N-2,N-1}, x_{N-1,N} + \sum_{i=2}^{N-3} k_{i,N-1}) B_{N-1}(x')', \end{aligned} \quad (4)$$

where

$$\begin{aligned} x_{ij} = -\alpha(s_{ij}) \quad \text{with} \quad s_{ij} = (p_i + p_{i+1} + \dots + p_j)^2, \\ z_{ij} = x_{ij} - x_{i+1,j} - x_{i,j-1} + x_{i+1,j-1}, \end{aligned}$$

and x'_{ij} are defined according to certain rules (given in a tabular form) which will not be reproduced here. More noteworthy of the present formula is the fact that Eq. (2) is self-contained such that in contrast to Ref. 3, there is required no prescription for defining the parameters of the function V_{n-1} which corresponds to B_{N-1} of Eq. (4).

In order to achieve the above we have to establish the relation between our parameters α_{ij} and those of Ref. 3. For this purpose let us note first that the external lines are labeled $1, 2, \dots, n+3$ both for the $(n+3)$ -point function in Ref. 1 and for V_n in the present paper, while they are labeled $0, 1, \dots, n+2$ for the $(n+3)$ -point function B_{n+1} of Ref. 7. Further, we note that the integration variables u_1, u_2, \dots, u_n in Ref. 7 and the present paper may be made to correspond to $u_{12}, u_{13}, \dots, u_{1,n+1}$ of Ref. 1. With this in mind one can compare Eq. (1) with the corresponding expression that follows from the representation for $B_{n+1}(x)$ of Ref. 1 through rearrangement of the integrand. Namely, by introducing $x_{ij} = -\bar{\alpha}_{ij} - 1$ from Ref. 1, where we write $\bar{\alpha}_{ij}$ for $\alpha_{ij} = \alpha(s_{ij})$ of Ref. 1, it becomes possible to express our α_{ij} in terms of $\bar{\alpha}_{km}$. If we further write $\xi_{ij} = -\bar{\alpha}_{ij}$ with $\xi_{ii} = 0$ and

$$\zeta_{ij} = \xi_{ij} - \xi_{i+1,j} - \xi_{i,j-1} + \xi_{i+1,j-1}, \quad (5)$$

where ξ_{ij} and ζ_{ij} stand for x_{ij} and z_{ij} , respectively, of Ref. 3, it follows that

AD753302

 DDC
 RECEIVED
 DEC 27 1972

$$\alpha_{0i} = \xi_{i,1,1} \quad \left\{ \begin{array}{l} \text{for } i = 1, 2, \dots, n \end{array} \right. \quad (6)$$

and

$$\alpha_{ki} = \xi_{i,1,1,k+1} \quad \text{for } k = 2, 3, \dots, n$$

$$\text{and } i = 1, 2, \dots, n - k + 1. \quad (7)$$

That the integral in Eq. (1) reduces for $w_{ij} = 1$ to B_{n+1} of Ref. 7 can be seen from the observation that our $p_i, i = 1, 2, \dots, n + 3$ correspond to $p_i, i = 0, 1, \dots, n + 2$ of Ref. 7 and through specialization of the relation $\bar{\alpha}(s_{ij}) = \alpha' s_{ij} + \alpha_j^0$ of Ref. 1 to the form $bs_{ij} + a$ as is done in Ref. 7.

Although we have connected our parameters to those of Ref. 3 in Eqs. (6) and (7), the precise correspondence between Eqs. (2) and (4) cannot be considered complete until the arrangement of ξ_{ij} in $B_N(\xi)$ of Eq. (4) is unambiguously established. [In Ref. 3 this arrangement has been left out unstated, which fact is responsible in part for requiring the somewhat troublesome rules for determining x' in $B_{N-1}(x')$ which should have really been unnecessary, as will be shown below.]

Let us suppose that the correspondence between V_n in this paper and B_N , for $N = n + 3$, of Ref. 3 is given by the following:

$$V_n(\alpha_{01}, \dots, \alpha_{0n}; \alpha_{11}, \dots, \alpha_{1n}; \alpha_{21}, \dots, \alpha_{2,n-1}; \dots; \\ \alpha_{n1}, w_{02}, \dots, w_{0n}; \dots; w_{n-2,n}) \\ \longleftrightarrow B_{n+3}(\xi_{12}, \xi_{23}, \dots, \xi_{n+1,n+2}; \xi_{13}, \xi_{24}, \dots, \\ \xi_{n,n+2}; \xi_{14}, \xi_{25}, \dots, \xi_{n-1,n+2}; \dots; \xi_{1,n+1}, \xi_{2,n+2}). \quad (8)$$

Then the transition from V_n to B_{n+3} and vice versa can be effected on a firm basis by referring to Eqs. (6) and (7).

With the help of Eq. (8) we now can translate Eq. (2) into a formula which is given in terms of the function B_N :

$$B_{n+3}(\xi_{12}, \xi_{23}, \dots, \xi_{n+1,n+2}; \xi_{13}, \xi_{24}, \dots, \\ \xi_{n,n+2}; \dots; \xi_{1,n+1}, \xi_{2,n+2})$$

$$= \sum \frac{(-\xi_{n,n+2}, r_{2,n+1}) \cdots (-\xi_{2,n+2}, r_{n+1})}{(1, r_{2,n+1}) \cdots (1, r_{n+1})} \\ \times B_4(\xi_{n+1,n+2}, \xi_{1,n+1} + \beta_{n,n+1}) \\ \times B_{n+2}(\xi_{12} + \beta_{n+1}, \xi_{23}, \dots, \xi_{n,n+1}; \xi_{13} + \beta_{n+2}, \\ \xi_{24}, \dots, \xi_{n-1,n+1}; \dots; \xi_{1n} + \beta_{n,n+1}, \xi_{2,n+1}). \quad (9)$$

Note that we wrote B_4 for B_4 and use was made of the following relation in obtaining Eq. (9):

$$B(\alpha_{0n}, \alpha_{1n}) (\alpha_{0n}, \beta_{n,n+1}) / (\alpha_{0n} + \alpha_{1n}, \beta_{n,n+1}) \\ = B(\alpha_{1n}, \alpha_{0n} + \beta_{n,n+1}).$$

That Eq. (9) is identical to Eq. (4) with $N = n + 3$, $x \rightarrow \xi$, and $z \rightarrow \xi$, can be checked easily. This establishes, therefore, that the order in which x_{ij} appears in $B_N(x)$ of Eq. (4), which was not stated explicitly in Ref. 3, should be exactly as is displayed in B_{n+3} of Eq. (8).

We emphasize that the recurrence formula for the $(n + 3)$ -point function, Eq. (9), as derived from Eq. (2) is complete as it stands and requires no rules for defining the parameters of the function B_{n+2} . In connection with the table for x'_{ij} of $B_{N-1}(x')$ in Ref. 3, we note that not all the entries in the table are actually needed for the recurrence formula. In fact, what is needed is only that portion of the table for $i = 1$, $j < N - 2$ and $i > 1$, $j \leq N - 2$ because, as may be seen from the arguments of B_{n+2} in Eq. (9), we require only ξ_{1j} for $j < N - 2 = n + 1$ and ξ_{ij} with $i > 1$ for $j \leq N - 2 = n + 1$. Moreover, there arises no need for including in the table the relation $x'_{1j} = x_{1j} + \sum_{k=2}^{j-1} k_{\alpha, N-1}$, for $j = N - 2$, unless we unnecessarily rewrite the argument $x_{N-1,N} + \sum_{k=2}^{N-3} k_{i, N-1}$ in B_4 of Eq. (4) as $x_{1,N-2} + \sum_{k=2}^{N-3} k_{i, N-1} = x'_{1,N-2}$. Finally, it is noted further that x'_{ij} for $i > 1$ and $j = N - 1$ should have not been included in the table since no such variables are actually involved in the recurrence formula for $B_N(x)$.

- 1 Chan Hong-Mo and Tsou Sheung Tsun, Phys. Letters 28B, 485 (1969).
- 2 C. J. Goebel and B. Sakita, Phys. Rev. Letters 22, 257 (1969).
- 3 J. F. L. Hopkinson and E. Plahte, Phys. Letters 28B, 489 (1969).
- 4 A. C. T. Wu, J. Math. Phys. 12, 2035 (1971).

- 5 P. Appell and J. Kampé de Fériet, Fonctions hypergéométriques et hypersphériques - polynômes d'Hermite (Gauthier-Villars, Paris, 1926), p. 114.
- 6 K. Mano, J. Math. Phys. 13, 1136 (1972).
- 7 K. Bardakci and H. Ruegg, Phys. Rev. 181, 1884 (1969).

